Numerical experiments in the problems of asymptotic representation theory

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Abstract

The article presents the results of experiments in computation of statistical values related to Young diagrams, including the estimates on maximum and average (by Plancherel distribution) dimension of irreducible representation of symmetric group S_n . The computed limit shapes of two-dimensional and three-dimensional diagrams distributed by Richardson statistics are presented as well.

1 Introduction

A classical definition of Young diagram[1] is the following: the Young diagram if size n is a finite descending ideal in the lattice $\mathbb{Z}_+ \times \mathbb{Z}_+$, i.e. a set of cells in positive quadrant, which contains a cell (i,j) only if it contains as well all the sells under (i,j) in terms of particular order.

Standard Young table of size n is a Young diagram whose cells are filled with numbers in the range from 1 to n, which increase in each row and each column. In other words, Young table is a path in a lattice of Young diagrams, which starts at the empty (zero-sized) diagram and ends in the given diagram.

Young diagrams correspond to irreducible representations of S_n (see for example [1, 2]), and the number of Young tables that fit into a given Young diagram is equal to the dimension of the corresponding irreducible representation. For the sake of brevity we call this number the dimension of the diagram Λ and denote it as $\dim(\Lambda)$.

This paper is organized as follows: in section 2 we will present the results of numerical experiments on asymptotics of a typical (by Plancherel measure) dimension of irreducible representation of the symmetric group. This measure was introduced in [3], for further details see [4].

Section 3 is devoted to the similar computations of a maximum dimension of irreducible representation of the symmetric group.

The results of section 2 should be considered as a support of a conjecture [4] about the existence of a limit value of normalized dimension of a typical Λ by Plancherel measure. This hypothetical limit value is denoted as c and called in [4] "specific entropy of an irreducible representation".

At the same time, basing on out computations in section 3, nothing conclusive can be said about the asymptotic behavior of the maximum dimension. That is, there is no claim that c_n comes to plateau on approachable values of n.

In section 4, we are concerned with another distribution on Young diagrams: Richardson distribution, which was studied in [5]. We give an experimental evidence of the proved theorem about a limit shape of typical Young diagrams. Then we make the similar experiments in three dimensions (this time the dimension has nothing to do with S_n), and form a hypothesis on a 3-dimensional limit shape.

2 Asymptotic behavior of a typical dimension of irreducible representation of S_n by Plancherel measure

Let \hat{S}_n be the collection of equivalence classes of complex irreducible representations of S_n . For $\Lambda_n \in \hat{S}_n$, we denote by dim Λ_n the dimension of the representation Λ_n , and by

$$\mu_n(\Lambda_n) = \frac{\dim^2 \Lambda_n}{n!} \tag{1}$$

we denote the Plancherel measure [3]. This is actually a probabilistic measure on \hat{S}_n , which we can derive from Burnside's conjecture:

$$\sum_{\Lambda_n \in \hat{S_n}} \dim^2 \Lambda_n = n!$$

The set \hat{S}_n and the dimension dim Λ_n have the following interpretation [1]: \hat{S}_n is the ensemble of all Young diagrams of size n, and dim Λ_n is the number of standard Young tables that fit into the diagram Λ_n . In the present article, we make no difference between Λ_n and the corresponding irreducible representation.

The following normalization of a dimension of diagrams with n cells (see [4]) is a key to studying the asymptotics of the dimension with $n \to \infty$:

$$c(\Lambda_n) = \frac{-2}{\sqrt{n}} \ln \frac{\dim \Lambda_n}{\sqrt{n!}}$$
 (2)

We will call the coefficient $c(\Lambda_n)$ the normalized dimension. From [4], the following two-sided estimates are known for $c(\Lambda_n)$:

$$\lim_{n \to \infty} \mu_n \left\{ \Lambda_n : c_0 < c(\Lambda_n) < c_1 \right\} = 1.$$

$$(c_0 = \frac{2}{\pi} - \frac{4}{\pi^2} \approx 0.2313, \ c_1 = \frac{2\pi}{\sqrt{6}} \approx 2.5651).$$

In other words, asymptotically almost all diagrams have the dimension lying inside the following range:

$$\sqrt{n!}e^{-\frac{c_1}{2}\sqrt{n}} < \dim \Lambda_n < \sqrt{n!}e^{-\frac{c_0}{2}\sqrt{n}}$$

Vershik and Kerov [4] stated the following conjecture: there exists a limit $\lim_{n\to\infty} c(\Lambda_n)$ almost everywhere (by Plancherel measure) on the set of infinite Young tables. Infinite Young table is an infinite sequence of nested Young diagrams of sizes increasing one-by-one: $\Lambda_1 \subset \Lambda_2 \subset \Lambda_3 \subset \ldots$ The hypothetical limit value is denoted as c and called "specific entropy of an irreducible representation".

We are studying the behavior of coefficient $c(\Lambda_n)$ w.r.t. n. Some experiments in this area were undertaken in [6], where the expectation and variance of $c(\Lambda_n)$ are given for five values of n, 1600 maximum (for n=1600 the sample was only 14 diagrams). We have done the same computations on modern hardware, which is far more powerful than 25 years ago.

Note. After we did the described experiments, it was reported of a new paper by Alexander Bufetov¹, which gives a proof of existence of a limit C of $c(\Lambda_n)$ in $L^2(Y)$ by Plancherel measure:

$$\lim_{n \to \infty} \int (c(\Lambda_n) - C)^2 d(\mu_n(\Lambda_n)) = 0$$

The actual value of C is not likely to be obtainable by methods used in Bufetov's proof. In should be also noted that the conjecture [4] about the existence of a limit if $c(\Lambda_n)$ almost everywhere is still open, and if it happens that such a limit value c exists, then clearly C = c.

The next two sections of the paper will describe well-known helper routines: generating the random diagrams by Plancherel measure with RSK algorithm and counting the dimension of the diagram.

2.1 Hook formula

Hook formula (see for example [1, 7]) allows to compute dim Λ , avoiding the enumeration of all Young tables that fit into Λ .

$$\dim \Lambda = \frac{n!}{\prod_{(i,j)\in\Lambda} h_Y(i,j)},\tag{3}$$

where (i, j) is a cell of Young diagram, and $h_Y(i, j)$ is the length of the hook associated to that cell. The hook of the cell (i, j) consists of the cell itself and all cells that are in j-th row to the right from (i, j) or in i-th column above (i, j) (fig. 1).

 $^{^1{\}rm Alexander}$ I. Bufetov. On the Vershik-Kerov Conjecture Concerning the Shannon-Macmillan-Breiman Theorem for the Plancherel Family of Measures on the Space of Young Diagrams. arXiv:1001.4275

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5	3	2	1	
7	5	4	3	1

Figure 1: Young diagram hook lengths

2.2 Generating RSK-random diagrams

While we have the hook formula in hand, the straight-forward generation of random diagrams according to Plancherel distribution would be very computationally expensive. The Robinson-Schensted-Knuth correspondence (RSK) and the row-insertion algorithm are of big help here.

The RSK algorithm [1] takes as input an arbitrary permutation $s \in S_n$, performs a sequence of row-insertions, and produces a pair of a standard Young tables (P,Q) whose diagrams are equal; furthermore, there is one-to-one correspondence (RSK) between such pairs and permutations. Hence, the uniform distribution (Haar measure) on S_n transforms into Plancherel measure on the set of left (or right) Young tableaux. Applying the RSK correspondence to a random permutation in S_n , and taking the Young diagram Y(P) of the left Young tableau in the pair, we will get a random Young tableau according to Plancherel measure.

2.3 Results

For $n \leq 120$ the expectancy of $c(\Lambda_n)$ for Plancherel measure were computed precisely by the formula,

$$c_n = \sum_{\Lambda_n \in \hat{S}_n} c(\Lambda_n) \frac{\dim^2 \Lambda_n}{n!}.$$

The results are listed in the next section: see table 3 and figure 5.

There are 1,844,349,560 Young diagrams of size 120. For larger n, the expectancies of $c(\Lambda_n)$ were calculated by Monte-Carlo method, using a sample of RSK-random diagrams. The normalized dimension of each diagram was computed by formula (2), using the hook formula 3 for obtaining dim Λ . The procedure was run against various n in the range from 1000 to 18000. The expectation value $c_n = E(c(\Lambda_n))$ and the standard deviation $\sigma_n = \sigma(c(\Lambda_n))$ of $c(\Lambda_n)$ are listed in table 1 and figure 2. From these results we can suggest that c_n is asymptotically increasing and allegedly has a limit. To be quite honest, we should notice that the c_n does not monotonically increase in the selected range: for example, it decreases from n = 15000 to n = 16000. This fact was re-checked and approved with a sample with 40000 items. The 3-rd decimal place remained constant after the amount of items had reached 20000.

n	sample size	$\approx c_n$	$\approx \sigma_n$
1000	2000	1.6984314	0.10431497
2000	2000	1.746588	0.091339454
3000	2000	1.7644972	0.08351989
4000	2000	1.7750576	0.07747431
5000	2000	1.7873781	0.07282907
6000	2000	1.7917556	0.07022077
7000	2000	1.7969893	0.06630529
8000	2000	1.8000197	0.06586118
9000	2000	1.8070668	0.06243244
10000	10000	1.8102994	0.061589677
11000	10000	1.8118591	0.059796795
12000	10000	1.8147597	0.057941828
13000	10000	1.8162445	0.05743194
14000	10000	1.8187699	0.056453623
15000	20000	1.820125	0.05504108
16000	20000	1.8181555	0.054255717
17000	20000	1.8197316	0.053651392
18000	20000	1.8249108	0.052745327

Table 1: Expected values and standard deviation of $c(\Lambda_n)$

2.4 Individual evolution of dimension of typical diagram

We define the Plancherel measure on infinite Young tableaux as a Markov measure having the following property: the corresponding measure on Young diagrams for each n is Plancherel measure (1) on diagrams. It is not a big problem to derive the transition probabilities for this Markov measure. Given that tables which have equal diagrams have as well equal measures, we can say that the measure of λ -shaped Young tableau is $\frac{\dim \lambda}{n!}$. Therefore, the probability of transition from λ to Λ is

$$P(\Lambda|\lambda) = \frac{\dim(\Lambda)}{(n+1)\dim(\lambda)}.$$
 (4)

(See for example [8]). So, the conjecture [4] on existence of limit of $c(\Lambda_n)$ almost everywhere means that for almost every infinite Young tableaux $\{\Lambda_n, n = 1, 2, ...\}$, generated by the described Markov process, $c(\Lambda_n)$ will converge to some common limit value, which obviously will be equal to the limit value for expectation of normalized dimension. Using the formula (4) for transition probability, we simulate the Markov process and obtain the sequence of Young diagrams of increasing size, each according to Plancherel measure, and each containing all the previous.

Our experiments have shown very chaotic behavior of the normalized dimension of such sequences, which probably points out that the Vershik-Kerov conjecture on existence of limit of $c(\lambda_n)$ almost everywhere is not easy to prove

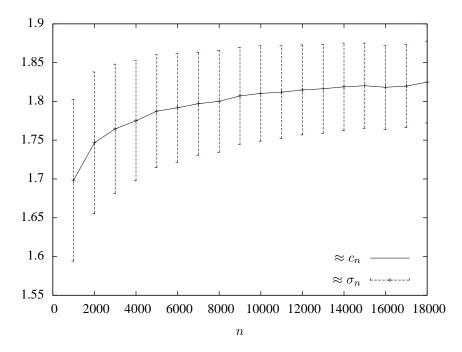


Figure 2: Expected values of $c(\Lambda_n)$. At each n, the height of a vertical line is equal to σ_n .

nor to prove wrong.

Non-regularity of behavior of normalized dimension is illustrated on figure 3, which depicts the values of $c(\lambda_n)$ for two Markov sequences of Young diagrams. The values are computed for $n \in [100..7000]$ (only multiples of 100 were taken).

3 Asymptotic behavior of a maximum dimension of irreducible representation of S_n

In this section we will study the behavior of $maximum\ dimension$ of a diagram of size n

$$m_n = \max_{\Lambda_n \in \hat{S_n}} \dim \Lambda_n$$

and its normalized value

$$\overline{c_n} = c(\overline{\Lambda}_n),$$

where $\overline{\Lambda}_n$ is the diagram of size n, which has a maximum dimension over all diagrams of size n.

The problem of computing maximum dimension was stated in 1968 (see [9]). In a paper by McKay [10] there are values of max dim Λ_n for n up to 75. Basing

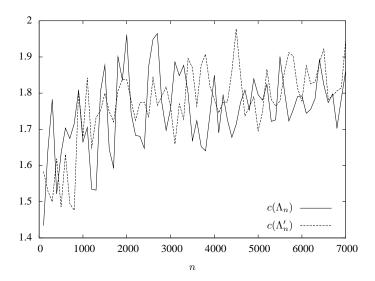


Figure 3: $c(\lambda_n)$ of two random sequences of Young diagrams by Plancherel measure

on his results, McKay assumed that

$$\max \frac{\dim \lambda_n}{\sqrt{n!}} \le \frac{1}{n}.\tag{5}$$

This assumption was the opposite to an alternative hypothesis, stating that there are irreducible representations of arbitrary large dimension, for which the inequality 5 is not true. Right before his paper was published, McKay sadly admitted that the alternative hypothesis is true for n=81. Nevertheless, as shown in [4], McKay's assumption is asymptotically true, and even stronger fact is true: $\frac{\max \dim \Lambda_n}{\sqrt{n!}}$ with $n \to \infty$ is decreasing as $e^{-c\sqrt{n}}$, i.e. not just faster than 1/n, but faster than any polynomial fraction. The estimates on normalized dimension, given in [4] for typical Young diagram, are the same for maximum dimension, and while both have the save logarithmic order, the experiment shows that the constants are different.

For n up to 130 we find the max dim Λ_n via enumeration of all Young diagrams of size n. In fact, we enumerated not diagrams, but partitions of n, using then the trivial correspondence between Young diagrams and partitions of integer [1]. The dimension of each diagram was computed using hook formula 3.

There are 5,371,315,400 Young diagrams of size n=130. For larger n, because of computational inability to enumerate all Young diagrams, the set of diagrams was restricted: we considered only symmetric diagrams, or diagrams that can be obtained from symmetric by adding one cell. This restriction often does not affect the final result, but for example for n=14 the diagram with maximum dimension does not pass. However, this restriction makes no big noise

after all. Table 2 contains values of $\overline{c_n}$. For $n=310,\,151,982,627$ diagrams were enumerated.

n	$\overline{c_n}$	n	$\approx \overline{c_n}$
10	0.57453286	140	1.05010306
20	0.8198125	150	1.0839802
30	0.7912792	160	1.05304872
40	0.86301332	170	1.0784368
50	0.90097636	180	1.0775954
60	0.94780416	190	1.0940416
70	0.98343194	200	1.0953336
80	0.96466594	210	1.1026434
90	0.9749938	220	1.11596048
100	1.035376	230	1.1106038
110	1.02168428	240	1.1273114
120	1.02246392	250	1.11251032
130	1.0514124	260	1.11878812
		270	1.1175388
		280	1.1173389
		290	1.13589692
		300	1.12641788
		310	1.148327

Table 2: Values of $c(\Lambda_n)$ for maximum-dimension diagrams. The first column has exact values, while the second column has the values obtained with enumeration of restricted set of diagrams.

While the two-sided estimates, given by Vershik and Kerov [4] for maximum dimension, are the same as for typical dimension, these two values have different behavior, and the limit of the sequence $\overline{c_n}$ is not likely to exist (see figure 4). We also emphasize that the maximum dimension is way greater than the typical one (vice versa for normalized values, because of the minus sign in the exponent).

Table 3 and figure 5 present the comparison of exact values of $\overline{c_n}$ and c_n . Neither of functions is monotonous.

4 Random diagrams by Richardson

Rost [5] considers a Markov process of a particle in $\{0,1\}^{\mathbb{Z}}$, which can be treated as a process of increasing the Young diagram cell-by-cell, starting from an empty diagram. The transition to the next state (increasing the diagram by one cell) is performed in the following way:

Among all diagrams of size n + 1, which contain the given diagram, one diagram is picked up randomly; each in the list with equal probability. In other words, from all the k "dimples" of the diagram of size n one dimple is picked

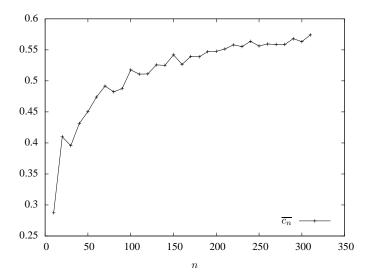


Figure 4: Values of $c(\Lambda_n)$ for maximum-dimension diagrams. The values starting from n = 140 are approximate, because of the mentioned restriction.

with probability 1/k. This growth process was introduced by Richardson [11]. Rost [5] found and proved the limit shape for Young diagram w.r.t. Richardson measure (see below).

We computed the values of normalized dimension $c(\Lambda_n)$ for sequences of nested Young diagrams, generated by Richardson process (see figure 7). The difference between these values and the values for Plancherel measure (figure 3) makes it clear that these measures are totally different.

Let us consider a process of infinite growth of Young diagram, with normalization on each step to along each axis; the area of the diagram thus remains constant. As shown in [5], with $n \to \infty$ the process converges to a limit shape, defined by equation $\sqrt{x} + \sqrt{y} = h$. The exact value of h is up to normalization. In [5] h is equal to one: $\sqrt{x} + \sqrt{y} = 1$. Another normalization that makes sense is the normalization by the area of the resulting figure, which is equal to 1/6 of the area of circumscribed square. The side of the square is equal to h^2 , so

$$S = \int_0^{h^2} (h - \sqrt{x})^2 dx = h^4/6$$

If we take the area S for 1, then the value of h is equal to $\sqrt[4]{6}$, and the limit-shape equation will be

$$\sqrt{x} + \sqrt{y} = \sqrt[4]{6} \tag{6}$$

n	$\overline{c_n}$	c_n
10	0.57453287	0.9348365
20	0.81981254	1.1238908
30	0.7912792	1.2205664
40	0.8630133	1.283057
50	0.90097636	1.3281072
60	0.94780415	1.3622344
70	0.98343194	1.3878295
80	0.96466595	1.4042087
90	0.9749938	1.4061089
100	1.035376	1.3848866
110	1.0216843	1.3299882
120	1.0224639	1.2363929

Table 3: Normalized dimensions: maximum and typical.

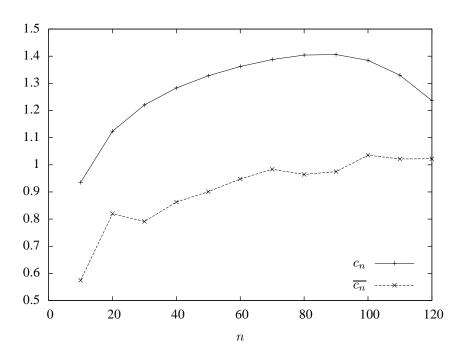


Figure 5: Normalized dimensions: maximum and typical.

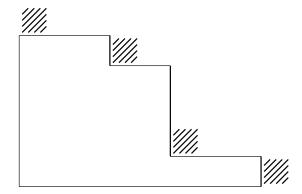


Figure 6: Building of a random Young diagram by Richardson measure. From k dashed dimples, each has the probability of 1/k.

4.1 d-dimensional Young diagrams

We call a d-dimensional Young diagram a finite descending ideal in the lattice $(\mathbb{Z}_+)^d$. Unless specified otherwise, just "Young diagram" will mean a two-dimensional Young diagram.

Vershik and Kerov [3] introduced a convenient coordinate system for presentation of Young diagrams: the diagram is rotated by 45° from the so-called French depiction, which conforms to Descartes coordinates.

Similarly, d-dimensional Young diagrams can be represented as functions, defined on a (d-1)-dimensional hyperplane, crossing the origin and orthogonal to the main diagonal. The value of the function is the length of a line span, parallel to the main diagonal, starting at the hyperplane and ending at the border of the Young diagram.

Having any Young diagram defined by this function, we easily define the average shape of a collection of diagrams, as the average of corresponding functions. This definition trivially applies to multi-dimensional Young diagrams as well.

Despite the function is defined in "rotated" coordinate system, we still depict the Young diagrams and their average shapes in Descartes coordinates, by applying an inverse transformation.

The average shape of 2200 random diagrams of size n = 100000 is shown on figure 8. A visual verification of Rost's theorem can be shown by plotting the average shape in coordinates (\sqrt{x}, \sqrt{y}) (see figure 9).

Scaling this shape by $1/\sqrt{n}$, we get the close approximation of the plot of equation 6. The area of the figure 10 is equal to 1.

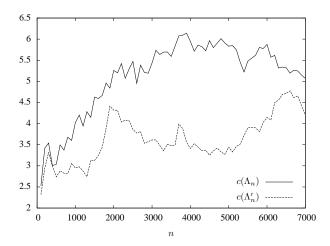


Figure 7: $c(\lambda_n)$ of two random sequences of Young diagrams by Richardson measure.

4.2 Standard deviation of main diagonal segment

In the previous section we verified that the average shape comply with equation 6. For verifying that the average shape is indeed a limit shape, we computed the standard deviation of s-called main diagonal segment—the length of a line span of the main diagonal, starting in origin and ending in the average shape. In table 4 the values of the standard deviation d(n) are listed, for n from 10000 to 40000, along with the normalized values $d(n)/\sqrt{n}$. Figure 11 shows the decrease of the normalized standard deviation with $n \to \infty$.

n	size of sample	$\approx d(n)$	$\approx d(n)/\sqrt{n}$
10000	2000	1.8262177	0.018262176
15000	2000	1.9742892	0.016120004
20000	3000	2.0621564	0.014581648
25000	4000	2.1949573	0.013882129
30000	5000	2.203268	0.012720575
35000	6000	2.3289392	0.012448704
40000	7000	2.3589768	0.011794884

Table 4: Standard deviation of main diagonal segment

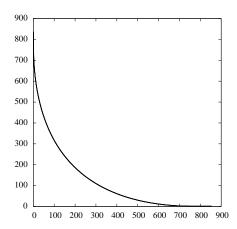


Figure 8: Average shape of 2200 diagrams of size 100000

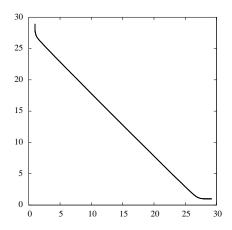


Figure 9: Average shape of 2200 diagrams of size 100000 in coordinates (\sqrt{x}, \sqrt{y})

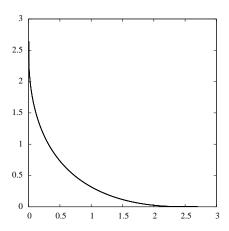


Figure 10: Average shape of 2200 diagrams of size 100000, scaled by $1/\sqrt{n}$

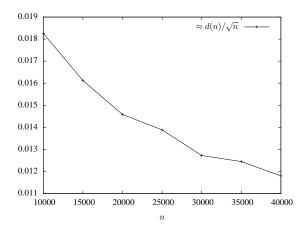


Figure 11: Normalized standard deviation of main diagonal segment

4.3 Average shape in three dimensions

The definition of random Young diagram by Richardson measure can be easily generalized to three-dimensional case. There are no known results about the limit shape in this case, but the figure 12, obtained by our computations, leads to assumption that the limit shape complies to similar equation $\sqrt{x} + \sqrt{y} + \sqrt{z} = h_3$.

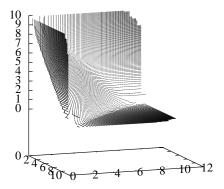


Figure 12: Average shape of 400 diagrams of size 10000, drawn in $(\sqrt{x}, \sqrt{y}, \sqrt{z})$ coordinates

This result shows that the limit shape of 3-dimensional Young diagrams generated by Richardson process are probably different from the limit shape for uniformly-distributed diagrams, which was studied in [12], and finally found in [13, 14].

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